

First Law Calculations for Ideal Gas Expansion and Compression

relationships that apply to **ideal gasses** for all conditions with $w_{\text{other}}=0$
(some also apply **more generally**):

$$\Delta U = q + w$$

$$w = -\int P_{\text{ext}} dV$$

$$PV = nRT$$

$$q_v = n\bar{C}_v \Delta T$$

$$q_p = n\bar{C}_p \Delta T$$

$$\bar{C}_p = \bar{C}_v + nR$$

$$H \equiv U + PV$$

$$\Delta U_{\text{any conditions}} = n\bar{C}_v \Delta T$$

$$\Delta H_{\text{any conditions}} = n\bar{C}_p \Delta T$$

monatomic ideal gas:

$$\bar{C}_v = \frac{3}{2}R$$

$$\bar{C}_p = \frac{5}{2}R$$

reversible vs irreversible processes:

Reversible: $P_{\text{ext}}=P_{\text{int}}=P$ throughout expansion or compression

Irreversible: often $P_{\text{ext}}=\text{constant}$

(external pressure instantaneously change to P_{ext}
and gas expands/contracts from P_{initial} to $P_{\text{final}}=P_{\text{ext}}$
doing work against $P_{\text{ext}}=\text{constant}$)

calculation of q , w , ΔU , ΔH for ideal gas:

1. any **isothermal** ($\Delta T=0$, $\Delta U=0$, $\Delta H=0$, $q=-w$)

a. *isothermal reversible*: $w = -\int P_{\text{ext}} dV = -\int_{V_1}^{V_2} \frac{nRT}{V} dV = -nRT \ln \frac{V_2}{V_1}$

b. *isothermal irreversible* (against $P_{\text{ext}}=\text{constant}$):

$$w = -\int P dV = -P_{\text{ext}} \int_{V_1}^{V_2} dV = -P_{\text{ext}} (V_2 - V_1)$$

2. any **adiabatic** ($q=0$, $w=\Delta U$)

a. *adiabatic reversible*: (here one would be given $(T_1, P_1, V_1) \rightarrow (P_2 \text{ or } V_2)$:
use relations among $P_1, T_1, V_1 \leftrightarrow P_2, T_2, V_2$ to get final T_2 :

$$\frac{T_1^{\frac{\bar{C}_p}{R}}}{P_1} = \frac{T_2^{\frac{\bar{C}_p}{R}}}{P_2}$$

$$P_1 V_1^{\frac{\bar{C}_p}{\bar{C}_v}} = P_2 V_2^{\frac{\bar{C}_p}{\bar{C}_v}}$$

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$$T_1^{\frac{\bar{C}_v}{R}} V_1 = T_2^{\frac{\bar{C}_v}{R}} V_2$$

for monatomic ideal gas: $\frac{\bar{C}_v}{R} = \frac{3}{2}$, $\frac{\bar{C}_p}{R} = \frac{5}{2}$, $\frac{\bar{C}_p}{\bar{C}_v} = \frac{5}{3}$

- ii once T_2 and thus ΔT are known then:

$$\Delta U = n\bar{C}_V\Delta T, \quad \Delta H = n\bar{C}_p\Delta T, \quad w = \Delta U, \quad q = 0$$

- b. *adiabatic irreversible* (against $P_{\text{ext}} = \text{constant}$):

- i here one would normally be given $(T_1, P_1, V_1) \rightarrow (P_2 = P_{\text{ext}})$ and one would need to find T_2 to apply relationships for ΔU and ΔH

$$w = - \int_{V_1}^{V_2} P_{\text{ext}} dV = -P_{\text{ext}}(V_2 - V_1) = P_{\text{ext}} \left(\frac{nRT_1}{P_1} - \frac{nRT_2}{P_2} \right)$$

$$\Delta U = n\bar{C}_V(T_2 - T_1) = w$$

$$n\bar{C}_V(T_2 - T_1) = P_{\text{ext}} \left(\frac{nRT_1}{P_1} - \frac{nRT_2}{P_2} \right)$$

with $P_2 = P_{\text{ext}}$ and solving for T_2 (factoring out n)

$$\left(\bar{C}_V + \frac{RP_{\text{ext}}}{P_1} \right) T_2 = \left(\bar{C}_V + \frac{RP_{\text{ext}}}{P_1} \right) T_1$$

$$T_2 = \left(\frac{\bar{C}_V + \frac{RP_{\text{ext}}}{P_1}}{\bar{C}_V + R} \right) T_1$$

you should verify if this makes sense for $P_{\text{ext}} > P_1$ (compression) versus $P_{\text{ext}} < P_1$ (expansion)

- ii once T_2 and thus ΔT are known then ;

$$\Delta U = n\bar{C}_V\Delta T, \quad \Delta H = n\bar{C}_p\Delta T, \quad w = \Delta U; \quad q = 0$$